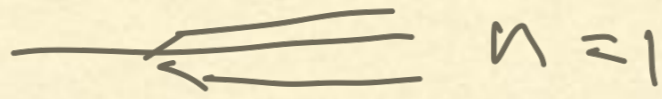


HI

continued



-  $\bar{t} = \frac{1}{n\gamma} = 350 \text{ yr}$  mean time b/w collisions

density

- spontaneous de-excitation

$$t = \frac{1}{A_{21 \text{ cm}}}$$

Einstein coefficient  
= measure of the probability  
of emission of light by  
atom or molecule  
A: spontaneous emission

= [table C.1] =

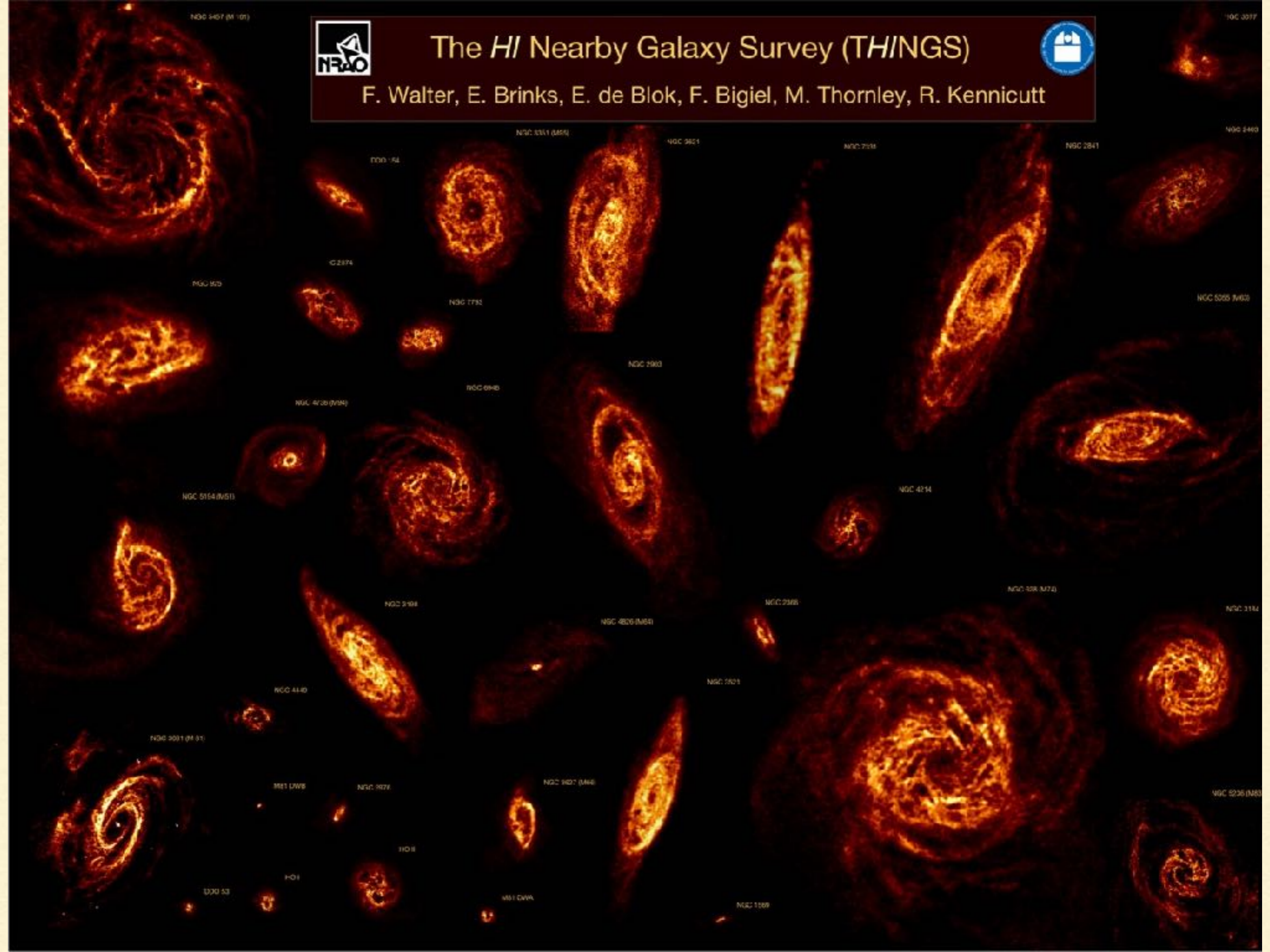
$$\frac{1}{2.976 \cdot 10^{-15} \text{ s}} = 10^7 \text{ yrs}$$



# The *HI* Nearby Galaxy Survey (*THINGS*)



F. Walter, E. Brinks, E. de Blok, F. Bigiel, M. Thornley, R. Kennicutt



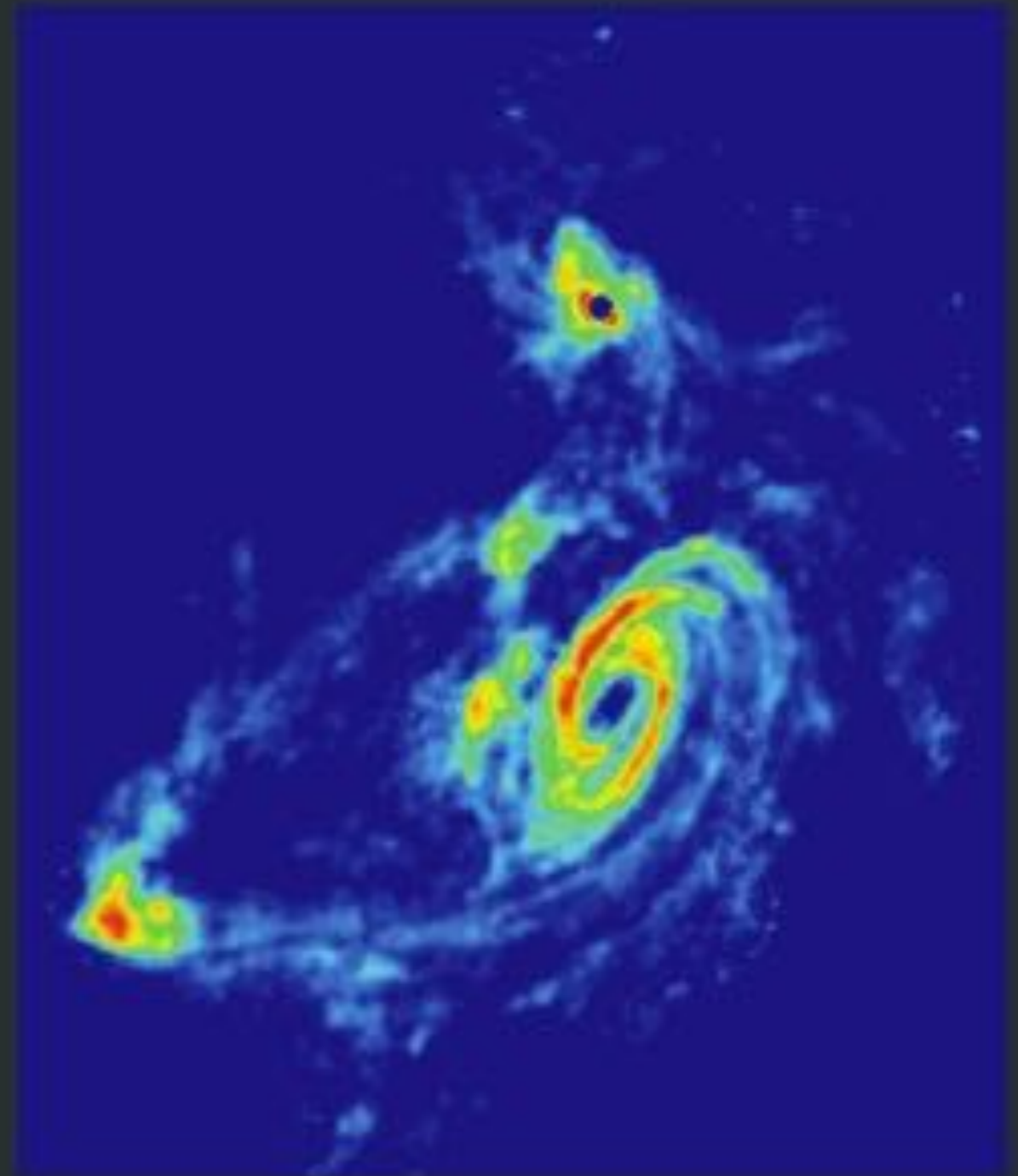


# TIDAL INTERACTIONS IN M81 GROUP

Stellar Light Distribution



21 cm HI Distribution



21 cm spectral line strong - why? Because there's so much HI! (App C  $\cup$ )

Ex) What  $\nabla$  is required (LTE) in order for equal # of particles in  $n=2$  and  $n=1$  states

eq 3.23 Boltzmann's equation

$$\frac{N_2}{N_1} = \frac{g_2}{g_1} e^{-\frac{\Delta E}{kT}} = \frac{n_2}{n_1} = 1$$

$$= \left[ g_n = 2n^2 \right] = \frac{8}{2} e^{-\frac{1.6 \cdot 10^{-11}}{kT}} = \left[ \begin{array}{l} \text{get } \Delta E \text{ from eq C6} \\ \Delta E = E_{n=1} - E_n = \end{array} \right.$$

$$= -13.6 \left( \frac{1}{n=2} - \frac{1}{n=1} \right) [\text{eV}] = -1.6 \cdot 10^{-11} \text{ eV} \quad \text{HOT!}$$

higher state \searrow lower state

Solve for  $T = T = 83634 \text{ K}$



$$T = 83000 \text{ K}$$

How many particles are ionized at this T?

eq 3.27

$$\frac{N_{\text{HII}}}{N_{\text{HI}}} = 2.41 \cdot 10^{15} \frac{T^{3/2}}{n_e} \cdot e^{-\frac{1.58 \cdot 10^5}{T}} = [T = 83634 \text{ K}]$$

$$= \frac{8,81 \cdot 10^{21}}{n_e}$$

$e^-$  density  $n_e$

for  $n_e \sim 10^{20} \text{ cm}^{-3}$

we require solar interior

Special case of  
Saha eq. 3.26.

# part  
in  $k+1$  state  
compared to  $k$   
state

So, H gas starts to become

appreciably ionized at T lower

than what's needed to excite the atom

(because more possible states available for a free  $e^-$   
than for a bound one in first excited state)

---

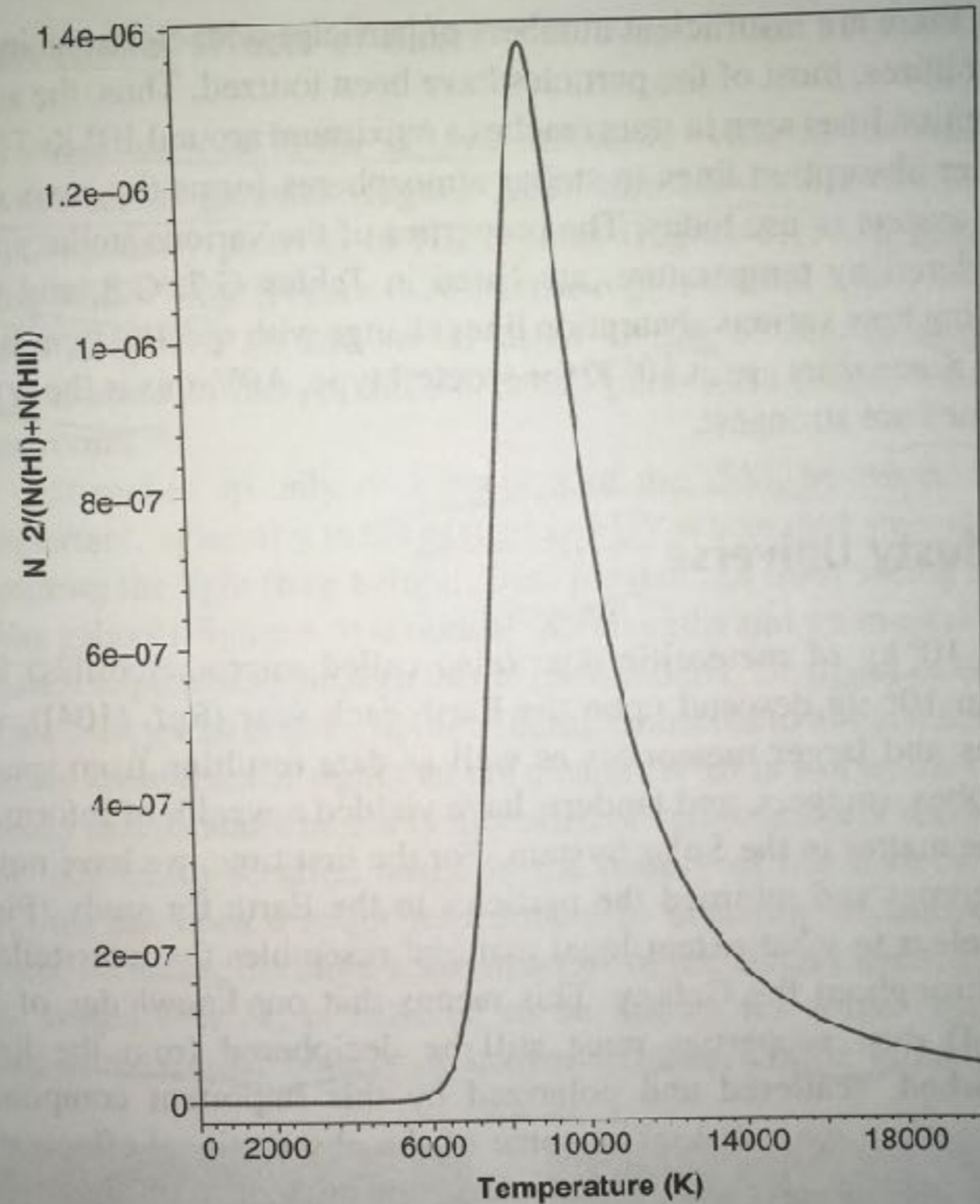
HI "all" particles are in ground state

How many particles are in state  $n=2$   
compared to all other states

$$\frac{N_2}{N_{\text{tot}}} = \frac{N_2}{N_{\text{HI}} + N_{\text{HII}}}$$



Balmer  
 absorption lines  
 $2 \rightarrow 3$   
 $2 \rightarrow 4$   
 around  
 $T = 8 - 10,000 \text{ K}$   
 highest # of  
 particles w/  
 $e^-$  in state  
 $n=2$



**Figure 3.18.** Fraction of hydrogen atoms that are in the first excited state,  $N_2$ , in comparison to the total number of hydrogen atoms,  $N_{\text{HI}} + N_{\text{HII}}$ . The adopted electron density at all temperatures has been taken to be constant at  $n_e = 10^{13} \text{ cm}^{-3}$  (in reality, this will also vary with temperature). This distribution indicates how the strength of the Balmer absorption lines varies with stellar surface temperature. Note that the fraction is still quite low, even at the peak.

3.5 Dust — extinction (dimming)  
                  — reddening (dust extinction is  
                  more effective @  
                  short  $\lambda$  than long  
                  ones  $\Rightarrow$  redder.)  
                  trigun stars & SNe

3.6 Cosmic rays — later. (synchrotron rad)



---

# CHAPTER 4: RADIATION ESSENTIALS

---

- Black body - Planck curve
- Rayleigh-Jeans & Wien laws

TE absorption & emission rates in balance  
T same!

⇒ radiation field isotropic and the  
spectrum depends only on T

Planck function / curve = Black body radiation

---

Planck function

$$B_{\nu}(T) = \frac{2h\nu^3}{c^2} \left( \frac{1}{e^{\frac{h\nu}{kT}} - 1} \right) \left[ \frac{\text{erg}}{\text{s cm}^2 \text{Hz sr}} \right] = T \quad 4.1$$

$$B_{\lambda}(T) = \frac{2hc^2}{\lambda^5} \left( \frac{1}{e^{\frac{hc}{\lambda T}} - 1} \right) \left[ \frac{\text{erg}}{\text{s cm}^2 \text{cm sr}} \right] = T \quad 4.2$$

Must be enough interaction in the gas w/ matter & radiation  $\Rightarrow$  must be opaque

$L <$  size of the object  $\leftrightarrow$  no reflection

BB is a perfect absorber  $\left\{ \begin{array}{l} \text{no reflection} \\ \text{no passing through} \end{array} \right.$

perfect emitter: abs & emission in balance

Single T

$=$  BB



reminder (abt  $I_\nu$  &  $I_\lambda$ ):

$$I_\nu d\nu = I_\lambda d\lambda$$

convert but need eq 1.5

$$B_\nu(T) d\nu = B_\lambda(T) d\lambda$$

$$d\lambda = \frac{-c}{\nu^2} d\nu$$

---

BB continuous Planck spectrum  
shape depends on temp

---

Perfect

BB:

CMBR.

despite

T gradient

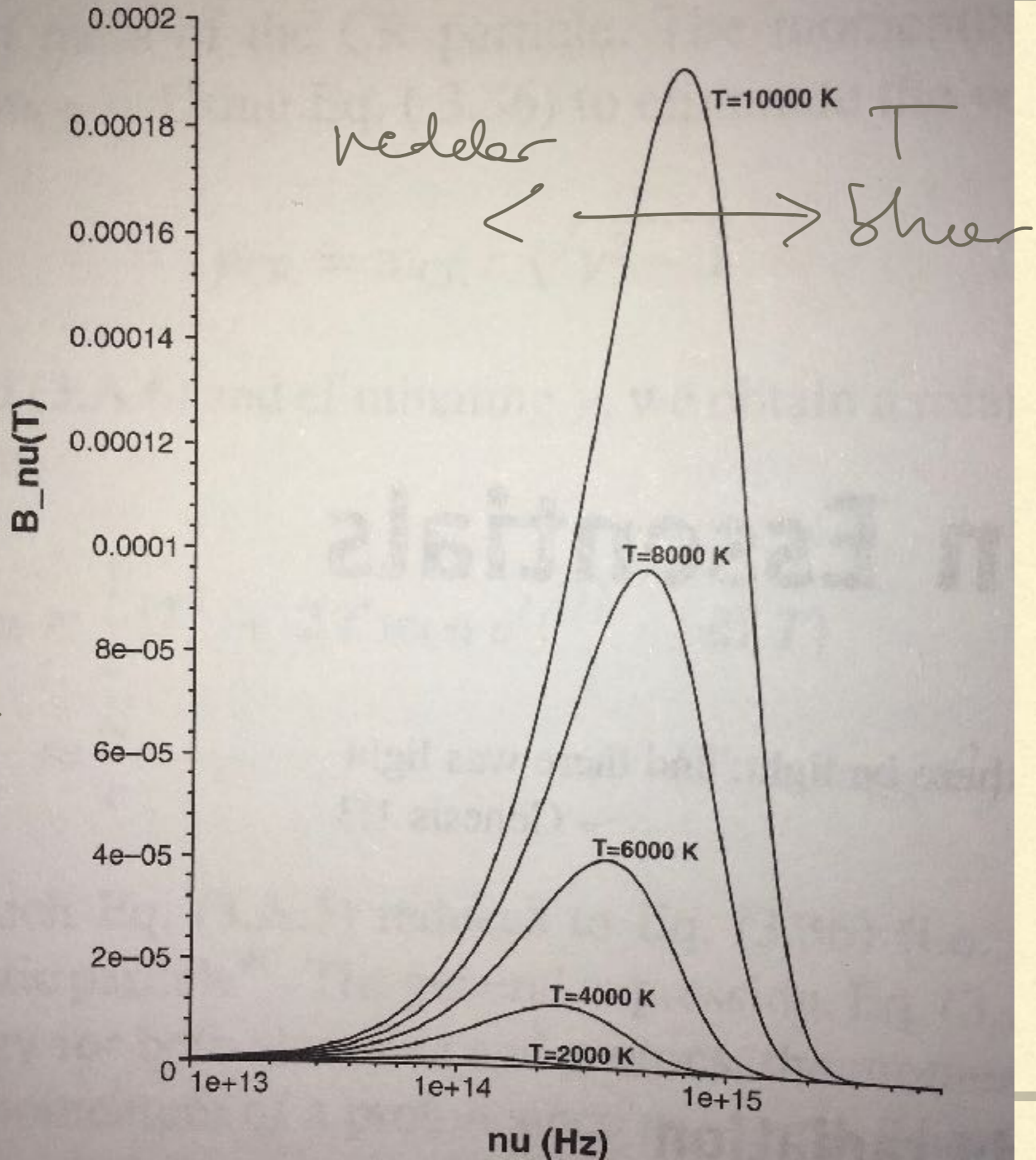
stars also

good BB

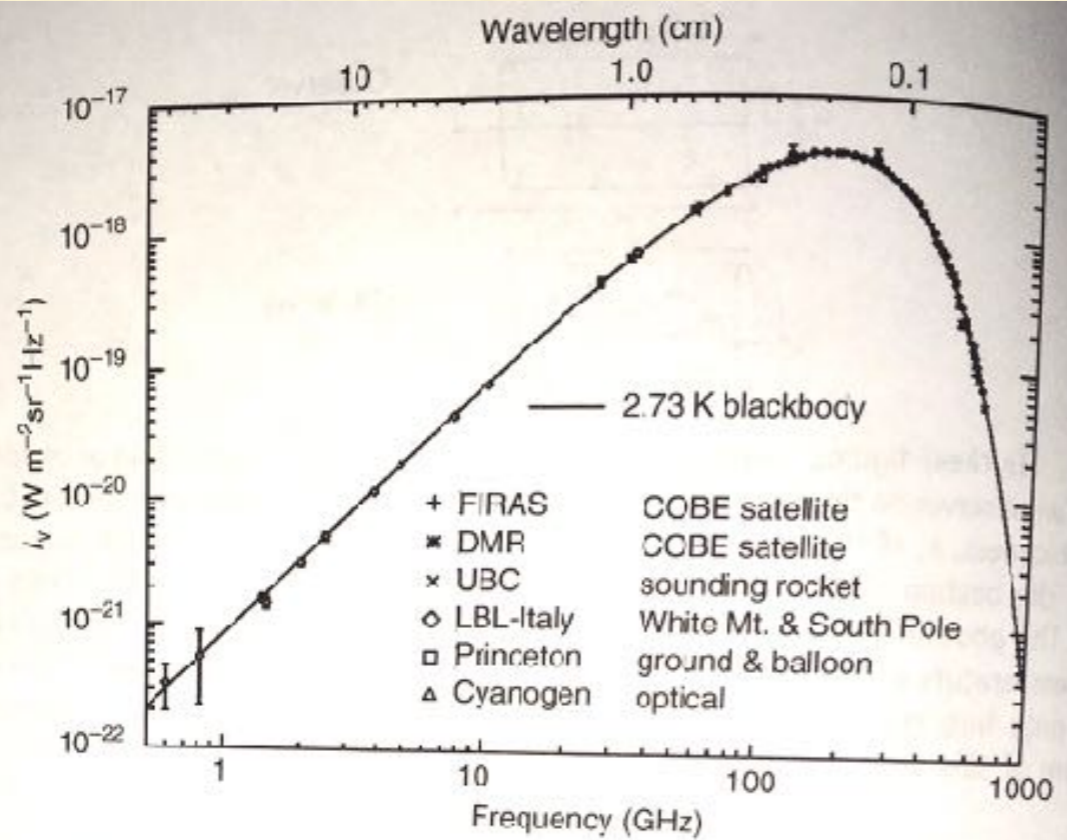
examples

with some

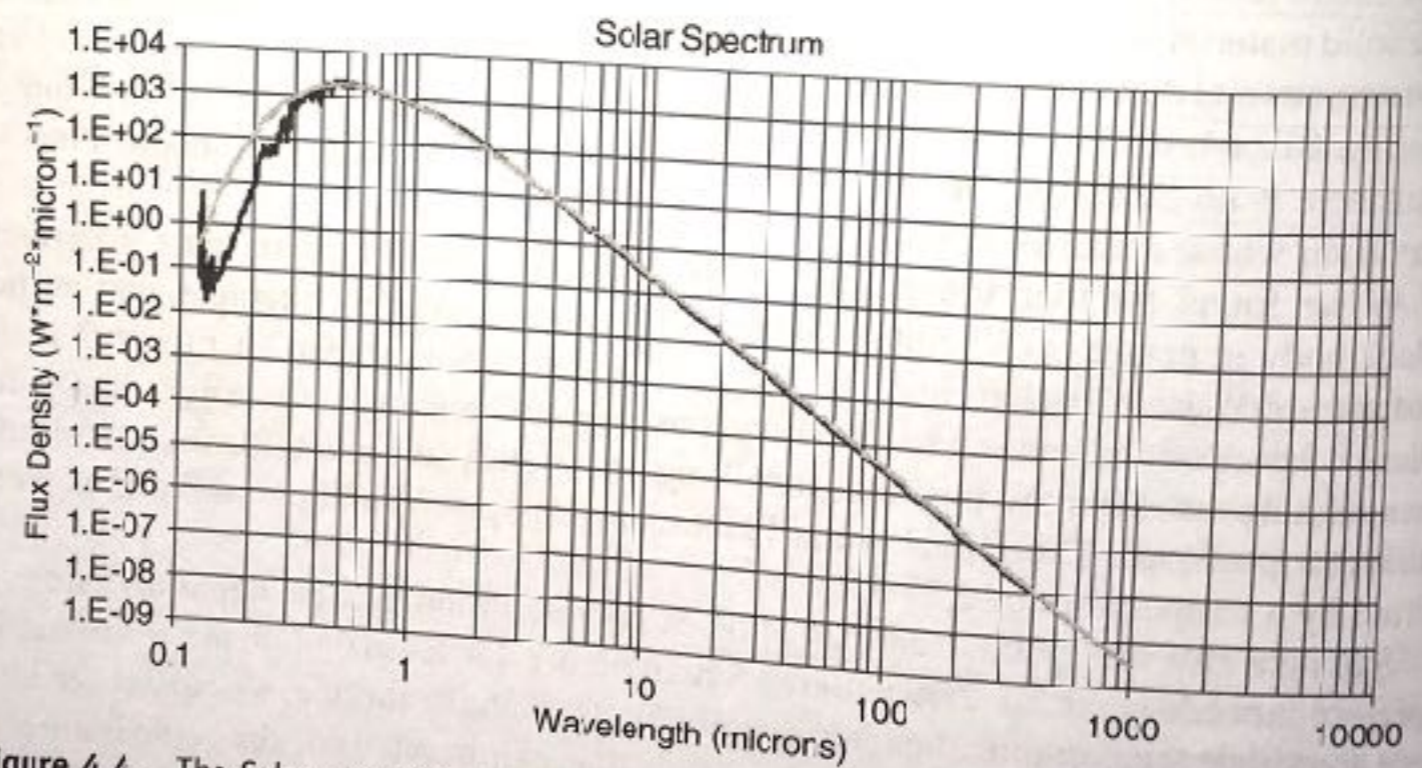
departures







**Figure 4.3.** Data points for this curve are measurements of the cosmic microwave background radiation from a variety of sources that are listed on the plot. The solid line is the blackbody spectrum, corresponding to a temperature of 2.73 K (G.F. Smoot et al., *Review of Particle Physics*, K. Hagiwara et al., Phys. Rev. **D66**, 010001, 1997).



**Figure 4.4.** The Solar spectrum, shown as a function of increasing wavelength in a logarithmic plot. The plot is of flux density as measured from Earth.